# MAT351 Partial Differential Equations Lecture 3 

September 21, 2020

## Last Lecture

## Theorem (Fundamental Theorem of Calculus of Variations, 1st version)

Consider $f \in C^{0}\left(\mathbb{R}^{n}\right)$. If

$$
\int f(\mathbf{x}) \varphi(\mathbf{x}) d \mathbf{x}=0 \quad \forall \varphi \in C_{c}^{0}\left(\mathbb{R}^{n}\right), \varphi \geq 0 \quad \Longrightarrow \quad f \equiv 0
$$

## Theorem (Fundamental Theorem of Calculus of Variations, 2nd version)

Consider $f \in C^{0}\left(\mathbb{R}^{n}\right)$. If

$$
\int_{\Omega} f(\mathbf{x}) d \mathbf{x}=\int f(\mathrm{x}) 1_{\Omega}(\mathrm{x}) d \mathrm{x}=0 \quad \forall \Omega \text { domain with smooth boundary in } \mathbb{R}^{n}
$$

$\Longrightarrow f \equiv 0$. A domain is connected and open subset of $\mathbb{R}^{n}$.

$$
1_{\Omega}(\mathbf{x})=\left\{\begin{array}{ll}
1 & \text { if } \mathbf{x} \in \Omega \\
0 & \text { if } \mathbf{x} \notin \Omega .
\end{array} \text { is the indicator or characteristic function of } \Omega \subset \mathbb{R}^{n} .\right.
$$

Hence, if $f, g \in C^{0}\left(\mathbb{R}^{n}\right)$ and

$$
\int_{\Omega} g(\mathrm{x}) d \mathrm{x}=\int_{\Omega} f(\mathrm{x}) d \mathrm{x} \forall \Omega \subset \mathbb{R}^{n} \text { with smooth boundary } \Rightarrow f \equiv g
$$

We derived 3 types of PDE:

- Simple Transport Equation:

$$
\begin{equation*}
u_{t}+V \cdot \nabla_{\mathbf{x}} u=0 \quad \text { in } \mathbb{R}^{n} \times \mathbb{R} \tag{1}
\end{equation*}
$$

- Scalar Conversation Laws:

$$
\begin{equation*}
u_{t}+\nabla_{\mathbf{x}} \mathbf{f}(u)=0 \quad \text { in } \mathbb{R}^{n} \times \mathbb{R} \tag{2}
\end{equation*}
$$

- Diffusion Equation:

$$
\begin{equation*}
u_{t}=-\lambda \Delta_{\mathrm{x}} u=-\nabla \cdot\left(\lambda \nabla_{\times} u\right) \quad \text { in } \mathbb{R}^{n} \times \mathbb{R} \tag{3}
\end{equation*}
$$

This equation is also known as heat equation since it also describes heat flow.
In this context the parameter $\lambda$ describes the heat conductivity, that also can depend on $\mathbf{x}$.
For simplicity, we will set $\lambda=1$ in the following.
If $u$ does not change in $t$, the left hand side is 0 . The equation becomes

- Laplace Equation

$$
0=\Delta u \quad \text { in } \mathbb{R}^{n}
$$

Solutions of this equation are called harmonic functions.

## General Solution

## A PDE

$$
\begin{equation*}
F\left(\mathbf{x}, u, D u, D^{2} u\right)=g(\mathbf{x}), \mathbf{x} \in \Omega \tag{4}
\end{equation*}
$$

can have infinitely many solutions.
The general solution is the collection of all $u$ that are sufficiently smooth and satisfy (4).

## Example

What is the general solution of the following PDE?

$$
\begin{equation*}
a u_{x}+b u_{y}=0, \text { in } \mathbb{R}^{2}, a \neq b . \tag{5}
\end{equation*}
$$

Rewrite (5) as $(a, b) \cdot \nabla u=0=\frac{\partial u}{\partial V}$ where $V=(a, b) \in \mathbb{R}^{2}$.
Hence, a solution $u$ must be constant in direction $V$.
The lines parallel to $V$ (Characteristics) have the form

$$
y=\frac{b}{a} x+\frac{c}{a} \text { or } b x-a y=c \text { where } c \in \mathbb{R} \text { is a parameter. }
$$

Therefore, the solution $u$ depends only on $b x-a y$. The general solution is

$$
u(x, y)=f(b x-a y) \text { for } f \in C^{1}(\mathbb{R})
$$

We can solve the previous PDE also using a coordinate change.
Define new coordinates via

$$
x^{\prime}(x, y)=a x+b y, \quad y^{\prime}(x, y)=b x-a y
$$

The chain rule yields

$$
u_{x}=u_{x^{\prime}} \frac{\partial x^{\prime}}{\partial x}+u_{y^{\prime}} \frac{\partial y^{\prime}}{\partial x}=a u_{x^{\prime}}+b u_{y^{\prime}}
$$

and similar $u_{y}=a u_{x^{\prime}}-b u_{y^{\prime}}$. Hence

$$
0=a u_{x}+b u_{y}=\left(a^{2}+b^{2}\right) u_{x^{\prime}} \Rightarrow u_{x^{\prime}}=0 \text { since } a^{2}+b^{2} \neq 0 .
$$

The general solution for this PDE is $f\left(y^{\prime}\right)=f(b x-a y)$ for $f \in C^{1}(\mathbb{R})$.

## Auxiliary Conditions

We want to specify an Auxiliary Condition that eventually yields a unique solution.
For instance, for

$$
a u_{x}+b u_{y}=0 \text { in } \mathbb{R}^{2}, a \neq b
$$

we could set the condition that $u(0, y)=g(y)$ for $g \in C^{1}(\mathbb{R})$.
The general solution is $f(b x-a y)=u(x, y)$ for $f \in C^{1}(\mathbb{R})$.
Let us determine $f$ such that $u$ satisfies the previous auxiliary condition:

$$
g(y)=u(0, y)=f(-a y)
$$

Hence, the solution is $u(x, y)=g\left(-\frac{1}{a}(b x-a y)\right)=g\left(y-\frac{b}{a} x\right)$.
Remark: Roughly, if a PDE has $n$ independent variable (in a domain $\Omega$ ) an auxiliary condition is a set of specified values on an $n-1$ dimensional subset of $\Omega$.

## Initial Value Conditions

Let $\phi \in C^{2}\left(\mathbb{R}^{n}\right)$.
We can search for solutions $u(\mathbf{x}, t)$ of the diffusion equation

$$
\begin{equation*}
u_{t}+\Delta u=0 \quad \text { on } \mathbb{R}^{n} \times\left[t_{0}, \infty\right) \tag{6}
\end{equation*}
$$

such that $u\left(\mathbf{x}, t_{0}\right)=\phi(\mathbf{x})$ for some $t_{0}$.

## Definition (Initial Value Problem (IVP) for the Diffusion Equation)

Let $\phi \in C^{2}\left(\mathbb{R}^{n}\right)$. The IVP for the diffusion equation is (6) together with

$$
\begin{equation*}
u(\mathbf{x}, 0)=\phi(\mathbf{x}) \text { for } \mathbf{x} \in \mathbb{R}^{n} \tag{7}
\end{equation*}
$$

## Definition (IVP for the Wave Equation)

The PDE $u_{t, t}=c^{2} \Delta u$ on $\mathbb{R}^{n} \times[0, \infty)$ is called Wave equation. The corresponding IVP is

$$
\begin{gathered}
u_{t, t}=c^{2} \Delta u \text { on } \mathbb{R}^{n} \times[0, \infty) \\
u(\mathbf{x}, 0)=\phi(\mathbf{x}) \text { and } u_{t}(\mathbf{x}, 0)=\psi(\mathbf{x}) \quad \text { for } \mathbf{x} \in \mathbb{R}^{n} .
\end{gathered}
$$

Here $c$ is a physical parameter that describes the speed of the wave.
We will study the wave equations later in more detail.

## Boundary Value Conditions

Let $\Omega$ be a domain with smooth boundary. Consider the Laplace equation

$$
\Delta u=0 \text { on } \Omega \text {. }
$$

The domain $\Omega \subset \mathbb{R}^{n}$ depends the (physical or mathematical) background.

## Definition

There are 3 common types of boundary conditions for the Laplace equation: Let $\phi, \psi \in C^{0}(\partial \Omega)$.

- Dirichlet boundary conditions (DC):

$$
u(\mathbf{x})=\phi(\mathbf{x}) \text { on } \partial \Omega .
$$

- Neumann boundary conditions (NC):

$$
\frac{\partial u}{\partial N}(\mathbf{x})=\phi \text { on } \partial \Omega
$$

where $\frac{\partial u}{\partial N}=\nabla u \cdot N$ is the derivative of $u$ in direction of $N$, and $N$ is the outer unit normal vector field of $\Omega$.

- Robin boundary conditions (RC):

$$
\frac{\partial u}{\partial N}(\mathbf{x})+\psi(\mathbf{x}) u(\mathbf{x})=\phi(\mathbf{x})
$$

If $\phi \equiv 0$ we call the corresponding boundary condition homogeneous, otherwise inhomogeneous.

## Example

The problem $\Delta u=0$ on $B_{1}(0)=\left\{|\mathbf{x}|_{2} \leq 1\right\}$ has infinitely many solutions $u \equiv c, c \in \mathbb{R}$.
Imposing a homogeneous DC picks the solution $u \equiv 0$.
On the other hand, any solution $u \equiv c, c \in \mathbb{R}$, satisfies the homogeneous NC. Hence, this problem is not wellposed.

We also can consider a general linear PDE of order 2:

$$
\sum_{i, j=1}^{n} a_{i, j} u_{x_{i}, x_{j}}+\sum_{k=1}^{n} b_{k} u_{x_{k}}+c u=0 \text { on } \Omega
$$

together with Dirichlet, Neumann or Robin boundary conditions.

## Example (Eigenvalue equation (in 1D))

Consider

$$
\frac{d}{d x} u+u=0 \text { on }[0, \pi] \subset \mathbb{R}
$$

A solution for the homogenenous Dirichlet problem (DP) is given by

$$
u(x)=c \sin x, \forall c \in \mathbb{R}
$$

a solution for the homogeneous Neumann problem (NP) is given by

$$
u(x)=c \cos x, \forall c \in \mathbb{R}
$$

## Global Condition/Condition at Infinity

One also can assume a global condition like

$$
\int_{\Omega} u(\mathbf{x}) d \mathbf{x}=0 \quad \text { or } \quad \int|u(\mathbf{x})|^{2} d \mathbf{x}=c \text { for } c>0 \text { fixed }
$$

## Example

The unique solution of the Neumann or Dirichlet eigenvalue problem such that

$$
\int_{0}^{\pi}|u(x)|^{2} d x=\pi
$$

is given by $u(x)=\cos x$ and $u(x)=\sin x$, respectively.
If $\Omega$ is unbounded (like $\Omega=\mathbb{R}^{n}$ or $\Omega=\mathbb{R}^{n-1} \times[0,1]$ ) one might also impose such an integral condition.
This forces $u$ to vanish for $|\mathbf{x}| \rightarrow \infty$.
Or we have directly a boundary condition of the form

$$
\lim _{|x| \rightarrow 0} u(\mathbf{x})=0 \text { or } \lim _{|x| \rightarrow 0} \frac{u(\mathbf{x})}{|\mathbf{x}|^{\alpha}}=0 \text { for } \alpha>0
$$

## Wellposed Problems

## Definition

Consider a PDE of order $k$

$$
F\left(\mathbf{x}, u, D u, \ldots, D^{k} u\right)=g(\mathbf{x})
$$

on a domain $\Omega \subset \mathbb{R}^{n}$ with a certain auxiliary condition.
The corresponding problem wellposed if it satisfies the following fundamental properties.
(1) Existence: There exists a solution.
(2) Uniqueness: The solution is unique.
(3) Stability: The unique solution depends in a stable manner on the data of the problem.

This means if we change the data or the auxiliary conditions a little bit then also the unique solution only changes a little bit.

## Remark

Since we often can measure data only up to certain degree of precision, stability guarantees that a solution of our problem that relies on the given data which approximates the exact data, is close to the solution that we would deduce from the exact data.

## Introduction to the Method of Characteristics

We found the general solution of $a u_{x}+b u_{y}=0$. Solutions are constant on lines parallel to $(a, b)$. Now, we consider

$$
u_{x}+y u_{y}=0 \Longleftrightarrow(1, y) \cdot \nabla u=0 \text { in } \mathbb{R}^{2} .
$$

Instead of straight lines, now we looking for curves $(x, y(x))$ such that

$$
\frac{d}{d x}(x, y(x))=(1, y) \Leftrightarrow \frac{d y}{d x}=y
$$

Hence $y(x)=C e^{x}$, and a solution $u$ satisfies

$$
\frac{d}{d x} u(x, y(x))=\nabla u \cdot(1, y)=0
$$

and

$$
u(x, y(x))=u(0, y(0))=u(0, C)
$$

is independent of $x$.
Let $g \in C^{1}(\mathbb{R})$.
Since for every tuple $(x, y)$ there exists a unique $C(x, y)$ such that $(x, y)=\left(x, C(x, y) e^{x}\right)$.
Then $u(x, y):=u(0, C(x, y))=g\left(y e^{-x}\right)$ satisfies

$$
u_{x}+y u_{y}=g^{\prime}\left(y e^{-x}\right)\left(-y e^{-x}\right)+y g^{\prime}\left(y e^{-x}\right) e^{-x}=0 .
$$

Therefore $u(x, y)=g\left(y e^{-x}\right)$ solves the PDE with auxiliary condition $g(y)=u(0, y)$.

