# MAT351 Partial Differential Equations Lecture 12 

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## Equations with a source term

In the following we will study inhomogeneous, linear second order PDEs
For instance, consider the initial value problem for diffusion equation with a source term:

## Diffusion equation with a source term

Let $f \in C^{0}(\mathbb{R} \times(0, \infty))$.

$$
\begin{array}{ccccc}
u_{t}-k u_{x, x} & =f(x, t) & \text { on } & \mathbb{R} \times(0, \infty) \\
u(x, 0) & =\phi(x) & \text { on } & \mathbb{R} .
\end{array}
$$

The physical interpretation of this equation is, for instance, the heat evolution of an infinitely long rod with an initial temperatur $\phi$ and a source (or sink) of heat at later times.

## Remark

If we define $A=k \frac{\partial^{2}}{\partial x^{2}}$, then $A$ is linear operator that goes from $C^{2}(\mathbb{R})$ to $C^{0}(\mathbb{R})$.
Then the inhomogeneous diffusion equation then takes the form

$$
\frac{d}{d t} u(t)=A u(t)+f(t) t>0 \text { and } u(0)=\phi \in C^{1}(\mathbb{R})
$$

where $u(t)=u(\cdot, t) \in C^{2}(\mathbb{R})$ and $f(t)=f(\cdot, t) \in C^{0}(\mathbb{R})$.

## Structural similarities with inhomogeneous ODEs

Recall the following ODE problem.

Let $A \in \mathbb{R}^{n \times n}$.

$$
\frac{d}{d t} v(t)=A v(t)+f(t), \quad v(0)=v_{0}
$$

where $t \in[0, \infty) \mapsto v(t), f(t) \in \mathbb{R}^{n}$.

For $f \equiv 0$ this is a homogeneous, linear ODE with constant coefficients.
The solution is given by $t \in[0, \infty) \mapsto e^{t A} v_{0}$.

$$
e^{t A} \text { is called the solution operator. }
$$

Recall: In case $A=B D B^{-1}$ for a diagonal matrix $D=\left(d_{1}, \ldots, d_{n}\right)$ then

$$
e^{t A}=B\left(e^{t d_{1}}, \ldots, e^{t d_{n}}\right) B^{-1}
$$

More general, one can find the operator $e^{t A}$ by means of the Jordan form for the matrix $A$.
The solution formula for the inhomogeneous problem with $f \neq 0$ is given by

$$
u(t)=e^{t A} v_{0}+\int_{0}^{t} e^{(t-s) A} f(s) d s
$$

## Dunhamel's principle

The solution formula for the inhomogeneous ODE is derived via Dunhamel's principle.
Assume $v(t)$ solves the inhomeogeneous problem. Assume $S(-t)=e^{-t A}=\left[e^{t A}\right]^{-1}$ exists.
Then we can compute

$$
S(-t) f(t)=S(-t)\left[\frac{d}{d t} v(t)-A v(t)\right]=S(-t) \frac{d}{d t} v(t)-S(-t) A v(t)=\frac{d}{d t}[S(-t) v(t)]
$$

The last equality is the product rule. Integrating from 0 to $t>0$ gives

$$
\int_{0}^{t} S(-s) f(s) d s=S(-t) v(t)-v_{0}
$$

Hence

$$
v(t)=S(t) v_{0}+S(t) \int_{0}^{t} S(-s) f(s) d s=S(t) v_{0}+\int_{0}^{t} S(t-s) f(s) d s
$$

We can then check that this $v(t)$ indeed satisfies the inhomogeneous ODE

$$
\begin{aligned}
\frac{d}{d t} v(t) & =\frac{d}{d t} S(t) v_{0}+\frac{d}{d t} \int_{0}^{t} S(t-s) f(s) d s \\
& =A S(t) v_{0}+S(0) f(t)+\int_{0}^{t} A S(t-s) f(s) d s \\
& =A\left[S(t) v_{0}+\int_{0}^{t} S(t-s) f(s) d s\right]+f(t)=A v(t)+f(t)
\end{aligned}
$$

The solution formula for ODEs gives us an idea how a solution formula for PDEs should look like.
We saw a version of this formula before in the case of inhomogeneous first order PDEs of the form

$$
u_{t}-a u_{x}=f(x, t)
$$

Here, the operator is given by $A=a \frac{\partial}{\partial x}$. Then the PDE takes the form

$$
u_{t}=A u+f(x, t)
$$

Recall the solution of the homogeneous equation was given by

$$
\phi(x+t a)=[S(t) \phi](x)
$$

Dunhamel's principle suggests the solution formula

$$
v(t)=S(t) \phi+\int_{0}^{t} S(t-s) f(s) d s=\phi(x+a t)+\int_{0}^{t} f(x+a(t-s), s) d s
$$

for the inhomogeneous problem.
This is exactly the formula that we already derived from the method of characteristics.

## Back to the inhomogeneous diffusion equation

The unique solution of the initial value problem

$$
\begin{array}{ccc}
u_{t} & =k u_{x, x} & \text { on } \mathbb{R} \times(0, \infty) \\
u(x, t) \rightarrow 0, & |x| \rightarrow \infty & \\
u(x, 0) & =\phi(x) & \text { on } \mathbb{R} .
\end{array}
$$

was given by

$$
\int_{-\infty}^{\infty} S(x-y, t) \phi(y) d y=: \mathbf{S}(t) \phi(x) \text { where } S(x, t)=\frac{1}{\sqrt{4 \pi k t}} e^{-\left(\frac{x}{\sqrt{4 k t}}\right)^{2}}
$$

We can see $\mathbf{S}(t): C^{1}(\mathbb{R}) \rightarrow C^{2}(\mathbb{R})$ as a family of solution operators.
Now we consider the same problem but with a source term $f \in C^{0}(\mathbb{R} \times(0, \infty))$ :

$$
u_{t}=k u_{x, x}+f(x, t) \quad \text { on } \mathbb{R} \times(0, \infty)
$$

We also write $f(s)$ for $f(\cdot, s)$. We assume $|f(x, t)| \leq C$. We prove the following theorem.

## Theorem

The unique solution ofthe inhomogeneous problem is given by the formula

$$
v(x, t)=[\mathbf{S}(t) \phi](x)+\int_{0}^{t}[\mathbf{S}(t-s) f(s)](x) d s
$$

Proof. We only check the existence statement.
First we compute $v_{t}=[\mathbf{S}(t) \phi]_{t}+\frac{d}{d t} \int_{0}^{t} \mathbf{S}(t-s) f(s) d s=k[\mathbf{S}(t) \phi]_{x, x}+\frac{d}{d t} \int_{0}^{t} \mathbf{S}(t-s) f(s) d s$. We consider the second term on the right hand side

$$
\frac{d}{d t} \int_{0}^{t}[\mathbf{S}(t-s) f(s)](x) d s=\frac{d}{d t} \int_{0}^{t} \int_{-\infty}^{\infty} S(x-y, t-s) f(y, s) d y d s=\frac{d}{d t} \int_{0}^{t} g(s, t) d s
$$

$s \in(-\infty, t] \mapsto g(s, t)$ is continuous with

$$
g(t, t)=\lim _{s \uparrow t} g(s, t)=\delta_{x}[f(\cdot, t)]=f(x, t)
$$

More precisely, we can compute

$$
\begin{aligned}
g(s, t) & =\int_{-\infty}^{\infty} S(x-y, t-s) f(y, s) d y \\
& =\int_{-\infty}^{\infty} S(x-y, t-s) f(y, t) d y+\int_{-\infty}^{\infty} S(x-y, t-s)(f(y, s)-f(y, t)) d y
\end{aligned}
$$

For the first term it follows by computation as we did before that

$$
\lim _{s \rightarrow t} \int_{-\infty}^{\infty} S(x-y, t-s) f(y, t) d y=\delta_{x}[f(t)]=f(x, t)
$$

For the second we get

$$
\begin{aligned}
\int_{-\infty}^{\infty} S(x-y)[f(y, s)-f(y, t)] d y= & \int_{x-\tilde{c}}^{x+\tilde{c}} S(x-y, t-s)[f(y, s)-f(y, t)] d y \\
& +\int_{\{y:|x-y|>\tilde{c}} S(x-y, t-s)[f(y, s)-f(y, t)] d y
\end{aligned}
$$

Since $f \in C^{0}(\mathbb{R} \times(0, \infty))$, $f$ is uniformily continuous on $[x-C, x+C] \times[t-\eta, t+\eta]$ for $t>0$ and $\eta>0$ sufficiently small such that $t-\eta>0$.
In particular, given $\epsilon>0$ there exists $\delta(\tilde{C}, \epsilon)>0$ such $|f(y, s)-f(y, t)|<\epsilon$ for $|s-t|<\delta$. Therefore, for the first term on the right hand side in the last formula we have

$$
-\epsilon \leq \int_{x-\tilde{C}}^{x+\tilde{C}} S(x-y, t-s)[f(y, s)-f(y, t)] d y \leq \epsilon
$$

For the second term on the right hand side in the last formula we have

$$
-2 C K e^{-\tilde{c}^{2}} \leq \int S(x-y, t-s)[f(y, s)-f(y, t)] \leq 2 C K e^{-\tilde{c}^{2}}
$$

because $|f(y, s)-f(y, t)| \leq C$ and $S(x-y, t-s) \leq K e^{-\tilde{C}^{2}}$ on $\{y:|x-y| \geq C\}$ for a constant $K>0$. So we can choose $\tilde{C}$ such that $2 C K e^{-\tilde{C}} \leq \epsilon$.

This considerations together imply that

$$
\lim _{s \rightarrow t} g(s, t)=f(x, t) \pm 2 \epsilon
$$

and since $\epsilon>0$ was arbitrary, the limit is $f(x, t)$.

Hence $\tau \in[0, \infty) \mapsto \int_{0}^{\tau} g(s, t) d s=\int_{0}^{\tau} \int_{-\infty}^{\infty} S(x-y, t-s) f(y, s) d y d s$ is differentiable in $t$ with

$$
\frac{d}{d t} \int_{0}^{\tau} g(s, t) d s=g(t, t)=f(x, t)
$$

Therefore

$$
\frac{d}{d t} \int_{0}^{t} g(s, t) d s=f(x, t)+\int_{0}^{t} \frac{\partial}{\partial t} g(s, t) d s
$$

For the second term on the right hand side we calculate

$$
\begin{gathered}
\int_{0}^{t} \frac{\partial}{\partial t} g(s, t) d t=\int_{0}^{t} \int_{-\infty}^{\infty} \frac{\partial}{\partial t} S(x-y, t-s) f(s) d y d s=\int_{0}^{t} \int_{-\infty}^{\infty} k \frac{\partial^{2}}{\partial x^{2}} S(x-y, t-s) f(s) d y d s \\
=k \frac{\partial^{2}}{\partial x^{2}} \int_{0}^{t} \int_{-\infty}^{\infty} S(x-y, t-s) f(s) d y d s=k\left[\int_{0}^{t}[\mathbf{S}(t-s) f(s)](x) d s\right]_{x, x}
\end{gathered}
$$

So we computed $v_{t}=f(x, t)+k\left[\mathbf{S}(t) \phi+\int_{0}^{t} \mathbf{S}(t-s) f(s) d s\right]_{x, x}$.

## Diffusion on the half line, Reflection method

We consider the Dirichlet problem for the diffusion equation:

$$
\begin{array}{ccc}
u_{t, t} & =k u_{x, x} & \text { on } \mathbb{R} \times(0, \infty) \\
u(x, 0) & =\phi(x) & \text { on }[0, \infty)  \tag{1}\\
v(0, t) & =0 & \text { for } t>0
\end{array}
$$

To find a solution formula for this equation we apply the reflection method.
Consider the odd extension of $\phi$ to the real line:

$$
\phi_{\text {odd }}(x)= \begin{cases}\phi(x) & x \geq 0 \\ -\phi(-x) & x<0\end{cases}
$$

The corresponding initial value problem has the solution

$$
u(x, t)=\int_{-\infty}^{\infty} S(x-y, t) \phi_{o d d}(y) d y
$$

Since $\phi_{\text {odd }}$ is odd, also $u(x, t)$ is odd, that is $u(x, t)=-u(-x, t)$ (Exercise). Hence $u(0, t)=0$ and the restriction $v$ of $u$ to $[0, \infty) \times[0, \infty)$ satisfies the Dirichlet boundary condition. Hence $v$ solves the Dirichlet problem for the diffusion equation with initial condition $\phi$. A formula of $v$ that only depends on $\phi$ (and not on $\phi_{\text {odd }}$ or $S(x, t)$ for $x \in \mathbb{R}$ ) is derived as follows

$$
\begin{gathered}
v(x, t)=\int_{0}^{\infty} S(x-y, t) \phi(y) d y+\int_{-\infty}^{0} S(x-y, t) \phi_{o d d}(y) d y \\
=\int_{0}^{\infty}[S(x-y, t) \phi(y)+S(x-y, t) \phi(-y)] d y=\int_{0}^{\infty}[S(x-y, t)-S(x+y, t)] \phi(y) d y
\end{gathered}
$$

The solution of the problem (1) is given by the formula

$$
v(x, t)=\int_{0}^{\infty}[S(x-y, t)-S(x+y, t)] \phi(y) d y
$$

Similar, we can consider the Neumann problem for the diffusion equation:

$$
\begin{aligned}
& u_{t} \quad=k u_{x, x} \quad \text { on }[0, \infty) \times(0, \infty) \\
& u(x, 0)=\phi(x) \quad \text { on } \mathbb{R} \\
& v_{x}(0, t)=0 \quad \text { for } t>0 .
\end{aligned}
$$

To derive a solution formula we apply the same strategy as for the Dirichlet problem: We consider the following initial value problem for the diffusion equation on the real line:

$$
\begin{array}{ccc}
u_{t, t} & =k u_{x, x} & \text { on } \mathbb{R} \times(0, \infty) \\
u(x, 0) & =\phi_{\text {even }}(x) & \text { on } \mathbb{R}
\end{array}
$$

where

$$
\phi_{\text {even }}(x)= \begin{cases}\phi(x) & x \geq 0 \\ \phi(-x) & x<0\end{cases}
$$

The solution of this initial value problem will be again even in $x: u(x, t)=u(-x, t)$.

## Diffusion with source term on the half line

Now we consider

$$
\begin{array}{ccc}
u_{t}-k u_{x, x} & =f(x, t) & \text { on }[0, \infty) \times(0, \infty) \\
u(x, 0) & =\phi(x) & \text { on } \mathbb{R}  \tag{2}\\
u(0, t) & =h(t) & \text { for } t>0
\end{array}
$$

for a boundary source function $h:[0, \infty) \rightarrow \mathbb{R}$ in $C^{1}([0, \infty))$.
A strategy to solve this problem is the Substraction method:
We consider $v(x, t)=u(x, t)-h(t)$. If $u \in C^{2}((0, \infty) \times(0, \infty))$ solves the previous problem, then $v \in C^{2}((0, \infty) \times(0, \infty))$ solves

$$
\begin{array}{cccc}
v_{t}-k v_{x, x} & = & f(x, t)-h^{\prime}(t) & \text { on }[0, \infty) \times(0, \infty) \\
v(x, 0) & = & \phi(x)-h(0) & \text { on } \mathbb{R} \\
u(0, t) & = & 0 & \text { for } t>0 .
\end{array}
$$

To solve this problem we can apply the reflection method as we did for the equation with $f \equiv 0$.
Then one can check that $v(x, t)+h(t)=: u(x, t)$ solves the problem (2).

