

MAT351 Partial Differential Equations

Lecture 14

November 2, 2020

Reflection method for wave equations

We will study the following Dirichlet problem for the wave equation on the half-line:

$$\begin{aligned}v_{t,t} &= c^2 v_{x,x} && \text{on } (0, \infty) \times \mathbb{R} \\v(x, 0) &= \phi(x) && \text{on } (0, \infty) \\v_t(x, 0) &= \psi(x) && \text{on } (0, \infty) \\v(0, t) &= 0 && \text{on } \mathbb{R}.\end{aligned}\tag{1}$$

The reflexion method works the same way as for the diffusion equation.

We consider *odd* extensions ϕ_{odd} and ψ_{odd} of ϕ and ψ respectively.

Let $u(x, t)$ be the solution of the initial value problem for the wave equation on \mathbb{R} . We have the formula

$$u(x, t) = \frac{1}{2} [\phi_{odd}(x + ct) + \phi_{odd}(x - ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi_{odd}(y) dy.$$

Then $u(x, t)$ is once again odd. In particular we have $u(0, t) = 0$ for $t > 0$ and we can define the solution v on $[0, \infty) \times \mathbb{R}$ of (1) by restriction of u to $[0, \infty)$.

We observe that for $x \geq c|t|$ it follows that $x - ct, x + ct \geq 0$. Hence

$$v(x, t) = \frac{1}{2} [\phi(x + ct) + \phi(x - ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(y) dy \quad x \geq c|t|.$$

For $0 < x < c|t|$ we have $\phi_{\text{odd}}(x - ct) = -\phi(-x + ct)$. Hence

$$v(x, t) = \frac{1}{2} [\phi(x + ct) + \phi(-x + ct)] + \frac{1}{2c} \int_{x-ct}^0 [-\psi(-y)] dy + \frac{1}{2c} \int_0^{x+ct} \psi(y) dy \quad 0 < x < c|t|.$$

We can apply a change of variable $y \mapsto -y$ to the first integral term. We obtain

$$\begin{aligned} v(x, t) &= \frac{1}{2} [\phi(ct + x) - \phi(ct - x)] + \frac{1}{2c} \int_{ct-x}^0 \psi(y) dy + \frac{1}{2c} \int_0^{x+ct} \psi(y) dy \\ &= \frac{1}{2} [\phi(ct + x) - \phi(ct - x)] + \frac{1}{2c} \int_{ct-x}^{ct+x} \psi(y) dy \quad 0 < x < c|t|. \end{aligned}$$

The complete solution is given by

$$v(x, t) = \begin{cases} \frac{1}{2} [\phi(x + ct) + \phi(x - ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(y) dy & \text{if } x \geq c|t| \\ \frac{1}{2} [\phi(ct + x) - \phi(ct - x)] + \frac{1}{2c} \int_{ct-x}^{ct+x} \psi(y) dy & \text{if } 0 < x < c|t|. \end{cases}$$

Finite Interval

Similarly we can also study the problem

$$\begin{aligned} v_{t,t} &= c^2 v_{x,x} && \text{on } (0, l) \times \mathbb{R} \\ v(x, 0) &= \phi(x) && \text{on } (0, l) \\ v_t(x, 0) &= \psi(x) && \text{on } (0, l) \\ v(0, t) = v(l, t) &= 0 && \text{on } \mathbb{R}. \end{aligned} \tag{2}$$

Diffusion equation with continuous initial data

Let us consider once more

$$\begin{aligned}u_t &= ku_{x,x} && \text{on } \mathbb{R} \times (0, \infty) \\ \lim_{t \downarrow 0} u(x, t) &= \phi(x) && \text{on } \mathbb{R}\end{aligned}$$

This time we assume $\phi \in C^0(\mathbb{R})$ and $|\phi(x)| \leq M \forall x \in \mathbb{R}$.

The convolution formula

$$u(x, t) = \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-\left(\frac{x-y}{\sqrt{4kt}}\right)^2} \phi(y) dy$$

still makes sense. Indeed, since $|\phi(x)| \leq M$ the integral is finite and bounded from above by M :

$$|u(x, t)| = \left| \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-\left(\frac{z}{\sqrt{4kt}}\right)^2} \phi(x-z) dz \right| \leq \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-\left(\frac{z}{\sqrt{4kt}}\right)^2} M dz \leq M.$$

A refined statement is that for $m \leq \phi(x) \leq M$ it follows

$$m \leq \int_{-\infty}^{\infty} S(x-y, t) \phi(y) dy \leq M \quad \forall t > 0 \quad (\text{Maximum Principle}).$$

Theorem

Let $\phi(x)$ and $u(x, t)$ be as above. Then $u \in C^\infty(\mathbb{R} \times (0, \infty))$ such that $u_t = ku_{x,x}$ on $\mathbb{R} \times (0, \infty)$ and $\lim_{t \downarrow 0} u(x, t) = \phi(x)$ for every $x \in \mathbb{R}$.

Proof of the theorem

We check that u is in $C^\infty(\mathbb{R} \times (0, \infty))$. Let $S(x, t) = \frac{1}{\sqrt{4\pi kt}} e^{-\left(\frac{x}{\sqrt{4\pi kt}}\right)^2}$.

We show that

$$\frac{\partial}{\partial x} \int S(x-y, t) \phi(y) dy = \int \frac{\partial}{\partial x} S(x-y, t) \phi(y) dy.$$

Recall that

$$\frac{\partial}{\partial x} S(x-y, t) \phi(y) = \lim_{h \rightarrow 0} \frac{1}{h} [S(x+h-y, t) - S(x-y, t)] \phi(y).$$

By the *dominated convergence theorem* for integrals we can pull this limit inside the integral if the modulus of the limit is bounded by an integrable function. This is indeed the case

$$\left| \frac{\partial}{\partial x} S(x-y, t) \phi(y) \right| \leq \left| -\frac{1}{\sqrt{4\pi kt}} \frac{x-y}{2kt} e^{-\frac{(x-y)^2}{4kt}} \right| M \leq \frac{M}{\sqrt{4\pi kt}} \frac{|x-y|}{2kt} e^{-\frac{|x-y|^2}{4kt}}.$$

The term on the right hand side has a finite integral on \mathbb{R} . Hence

$$\frac{\partial}{\partial x} u(x, t) = \frac{1}{\sqrt{4k\pi t}} \int_{-\infty}^{\infty} \frac{\partial}{\partial x} S(x-y, t) \phi(y) dy.$$

All other derivatives of higher order in x and t will work the same way: we always get an estimate by function of the form

$$C|y-x|^n e^{-\tilde{C}(x-y)^2}$$

that has finite integral on \mathbb{R} .

Checking the initial condition

We also know that u satisfies $u_t = ku_{x,x}$ because $S(x, t)$ does.

Hence, we only need to prove that u satisfies the initial condition for $t \downarrow 0$.

Consider

$$u(x, t) - \phi(x) = \int_{-\infty}^{\infty} S(x-y, t)\phi(y)dy - \int_{-\infty}^{\infty} S(x-y, t)\phi(x)dy = \int_{-\infty}^{\infty} S(x-y, t)(\phi(y) - \phi(x))$$

Since ϕ is continuous in x , for $\epsilon > 0$ we can choose $\delta > 0$ such that

$$|y - x| \leq \delta \Rightarrow |\phi(x) - \phi(y)| \leq \epsilon$$

Hence

$$\begin{aligned} |u(x, t) - \phi(x)| &\leq \int_{\{y \in \mathbb{R}: |x-y| > \delta\}} \frac{1}{\sqrt{4\pi kt}} e^{-\frac{(x-y)^2}{2kt}} \underbrace{|\phi(x) - \phi(y)|}_{\leq 2M} dy \\ &\quad + \int_{\{y \in \mathbb{R}: |x-y| \leq \delta\}} S(x-y, t) \underbrace{|\phi(x) - \phi(y)|}_{\leq \epsilon} dy \\ &\leq \frac{2M}{\sqrt{4\pi}} \int_{\{z \in \mathbb{R}: |z| \geq \frac{\delta}{\sqrt{kt}}\}} e^{-\frac{z^2}{4}} dz + \epsilon. \end{aligned}$$

It follows that

$$\limsup_{t \downarrow 0} |u(x, t) - \phi(x)| \leq \epsilon$$

Since $\epsilon > 0$ was arbitrary, it follows that $\lim_{t \downarrow 0} |u(x, t) - \phi(x)| = 0$. □

Additional Remarks

- Decay of the solution for $t \rightarrow \infty$.
For $\phi \in C^0(\mathbb{R})$ with $|\phi| \leq M$ we have

$$|u(x, t)| \leq \int_{-\infty}^{\infty} S(x-y, t) |\phi(y)| dy \leq \frac{M}{\sqrt{4k\pi t}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4kt}} dy \leq \frac{M}{\sqrt{4\pi kt}} \rightarrow 0.$$

In particular, this means the backwards diffusion equation

$$u_t = -ku_{x,x} \text{ on } \mathbb{R} \times (0, \infty)$$

is not well-posed because stability fails.

- About uniqueness again: Let $\phi_1, \phi_2 \in C^0(\mathbb{R})$ with $|\phi_1|, |\phi_2| \leq M$.
We saw that in the class of solutions with $u(x, t) \rightarrow 0$ for $|x| \rightarrow \infty$ we find a unique solution.
But if we drop this assumption uniqueness might fail: There are solutions of the heat equation with $u(x, t) \rightarrow 0$ for $t \downarrow 0$ for all $x \in \mathbb{R}$.
See also exercise 10 on page 399 in Choksi's Lecture Notes for an example that hints to nonuniqueness.