

MAT351 Partial Differential Equations

Lecture 19

November 30, 2020

Last Lecture

Theorem

Let $f \in C^1(\mathbb{R})$ be periodic with period 2π .

Then the full Fourier series of f

$$\frac{1}{2}\tilde{A}_0 + \sum_{n=1}^{\infty} [A_n \sin(nx) + \tilde{A}_n \cos(nx)]$$

converges on $[-\pi, \pi]$ uniformly to f . Recall that the coefficients are given by

$$A_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin(nx) dx, \quad \tilde{A}_0 = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) dx, \quad \tilde{A}_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \cos(nx) dx.$$

We also showed that

$$\frac{1}{2}|\tilde{A}_0| + \sum_{n=1}^{\infty} (|A_n| + |\tilde{A}_n|)$$

Note that

$$A_n = \frac{1}{\|X_n\|^2} (f, X_n), \quad \tilde{A}_n = \frac{1}{\|\tilde{X}_n\|^2} (f, \tilde{X}_n) \quad \text{for } n \in \mathbb{N}$$

where $X_n(x) = \sin(nx)$, $\tilde{X}_n(x) = \cos(nx)$ and $(f, g) = \int_{-\pi}^{\pi} f(x)g(x)dx$ for $f, g \in C^0(\mathbb{R})$ periodic.

But for $\tilde{X}_0(x) = 1$ the definition says

Fourier sine and cosine series

Consider $f \in C^1([0, \pi])$ with Dirichlet boundary conditions: $f(0) = f(\pi) = 0$.

Let f_{odd} be the odd periodic extension of f to \mathbb{R} . Then $f_{\text{odd}} \in C^1(\mathbb{R})$ because

$$\lim_{h \downarrow 0} \frac{f(h) - f(0)}{h} = \lim_{h \downarrow 0} \frac{f(h)}{h} = \lim_{h \downarrow 0} \frac{-f(h)}{-h} = \lim_{h \uparrow 0} \frac{f_{\text{odd}}(h)}{h}.$$

We know that the full Fourier series of f_{odd} has the form

$$\sum_{n=1}^{\infty} A_n \sin(nx) = \mathcal{F}(f)$$

and converges uniformly on $[-\pi, \pi]$ to f_{odd} .

Recall that

$$\frac{1}{\pi} \int_{-\pi}^{\pi} f_{\text{odd}}(x) \sin(nx) dx = \frac{1}{\pi} \int_{-\pi}^0 \underbrace{f_{\text{odd}}(x)}_{-f(-x)} \sin(nx) dx + \frac{1}{\pi} \int_0^{\pi} f(x) \sin(nx) dx = \frac{2}{\pi} \int_0^{\pi} f(x) \sin(nx) dx$$

Hence $A_n = A_n^{\text{sin}}$ where A_n^{sin} is the Fourier sine coefficient.

Therefore, the full Fourier series $\mathcal{F}(f_{\text{odd}})$ coincides with the Fourier sine series $S(f)$ of f on $[0, \pi]$ and we have the following

Corollary

Let $f \in C^1([0, \pi])$ with $f(0) = f(\pi) = 0$.

Then, the Fourier sine series converges uniformly to f on $[0, \pi]$.

For $f \in C^1([0, \pi])$ with Neumann boundary conditions ($f'(0) = f'(\pi) = 0$) we consider the even periodic extension f_{even} .

Then again $f_{\text{even}} \in C^1(\mathbb{R})$, the full Fourier series converges uniformly to f and $\tilde{A}_n = A_n^{\text{cos}}$ where A_n^{cos} are the coefficients of the Fourier cosine series.

We obtain

Corollary

For $f \in C^1([0, \pi])$ with Neumann boundary conditions, the Fourier cosine series of f converges uniformly to f on $[0, \pi]$.

Recall that more generally one has

Theorem

Let $f \in C^0(\mathbb{R})$ be periodic with period 2π .

Then the full Fourier series of f converges uniformly on $[-\pi, \pi]$ to f .

Corollary

Let $f \in C^0([0, \pi])$ with $f(0) = f(\pi) = 0$. Then, the Fourier sine series converges uniformly to f on $[0, \pi]$.

For $f \in C^1([0, \pi])$ with Neumann boundary conditions, the Fourier cosine series of f converges uniformly to f on $[0, \pi]$.

Application: Heat equation with Dirichlet boundary conditions on $[0, \pi]$

Let $f \in C^1([0, \pi])$ with Dirichlet boundary conditions and let f_{odd} be the odd periodic extension. Then f has an expansion as Fourier sine series:

$$\sum_{n=1}^{\infty} A_n \sin(nx) = f(x).$$

Theorem

The series

$$u(x, t) = \sum_{n=1}^{\infty} A_n e^{-kn^2 t} \sin(nx)$$

converges uniformly on $[0, \pi] \times [0, \infty)$.

We have $u \in C^2([0, \pi] \times (0, \infty)) \cap C^0([0, \pi] \times [0, \infty))$ and u solves

$$u_t = ku_{x,x} \quad \text{on } [0, \pi] \times (0, \infty)$$

$$u(x, 0) = f(x) \quad \text{on } [0, \pi]$$

$$u(0, t) = 0 \quad \text{and} \quad u(\pi, t) = 0 \quad \forall t > 0.$$

Proof of the theorem

First, we have for $N, M \in \mathbb{N}$ and $M > N$

$$\begin{aligned} & \max_{(x,t) \in [0,\pi] \times [0,\infty)} \left| \sum_{n=1}^M A_n e^{-ktn^2} \sin(nx) - \sum_{n=1}^N A_n e^{-kn^2 t} \sin(nx) \right| \\ &= \max_{(x,t) \in [0,\pi] \times [0,\infty)} \left| \sum_{n=N+1}^M A_n e^{-kn^2 t} \sin(nx) \right| \\ &\leq \sum_{n=N+1}^M |A_n| e^{-kn^2 t} |\sin(nx)| \rightarrow 0. \end{aligned}$$

We use that $\sum_{n=1}^{\infty} |A_n| < \infty$ is finite, what was part of the proof of uniform convergence of the Fourier series for $f \in C^1(\mathbb{R})$ periodic.

Therefore, the partial sums $\sum_{n=1}^N A_n e^{-kn^2 t} \sin(nx)$ are a Cauchy sequence w.r.t. to uniform convergence, and hence, the uniform limit

$$u(x, t) = \sum_{n=1}^{\infty} A_n e^{-ktn^2} \sin(nx)$$

exists and $u(x, t)$ is a continuous function on $[0, \pi] \times [0, \infty)$ and $u(x, t)$ satisfies $u(x, 0) = f(x)$, $u(0, t) = u(\pi, t) = 0$. In particular, u is continuous on $[0, \pi] \times [0, \infty)$.

Moreover, each term $A_n e^{-ktn^2} \sin(nx)$, $n \in \mathbb{N}$, has partial derivatives w.r.t. t and x :

$$-A_n n^2 k e^{-ktn^2} \sin(nx) \quad \text{and} \quad A_n n e^{-ktn^2} \cos(nx)$$

and second derivatives w.r.t. x : $-A_n n^2 e^{-ktn^2} \sin(nx)$ and solves the heat equation by the separation of variable method. Recall the series $\sum_{n=1}^{\infty} n^\alpha e^{-ktn^2}$ converges absolutely whenever $t > 0$ for all $\alpha \in \mathbb{N}$. Hence, it follows that the series

$$\sum_{n=1}^{\infty} A_n n e^{-ktn^2} \sin(nx), \quad \sum_{n=1}^{\infty} A_n n^2 e^{-ktn^2} \sin(nx)$$

converge uniformly on $[0, \pi] \times [t_0, \infty)$ for $t_0 > 0$. This follows because, for instance,

$$\begin{aligned} \max_{(x,t) \in [0,\pi] \times [t_0,\infty)} \left| \sum_{n=N+1}^{\infty} A_n n e^{-ktn^2} \sin(nx) \right| &\leq \sum_{n=N+1}^{\infty} |A_n| n e^{-kt_0 n^2} \\ \text{(Cauchy-Schwartz inequality)} &\leq \sqrt{\left(\sum_{n=N+1}^{\infty} A_n^2 \right) \left(\sum_{n=N+1}^{\infty} n^2 e^{-2kt_0 n^2} \right)} \end{aligned}$$

It follows that we can compute the first and second order derivatives of $u(x, t)$ w.r.t. x and t by computing the partial derivatives of the partial sums:

$$u_x(x, t) = \sum_{n=1}^{\infty} A_n e^{-ktn^2} n \cos(nx), \quad u_t(x, t) = \sum_{n=1}^{\infty} A_n k n^2 e^{-ktn^2} \sin(nx)$$

on $[0, \pi] \times [t_0, \infty)$ for $t_0 > 0$ and

$$k u_{x,x}(x, t) = -k \sum_{n=1}^{\infty} A_n n^2 e^{-ktn^2} \sin(nx) = u_t(x, t) \text{ for } (x, t) \in [0, \pi] \times (0, \infty).$$

In particular, $u \in C^2([0, \pi] \times (0, \infty))$ and solves the heat equation. □

Complex form of the full Fourier series

Recall that

$$\sin(nx) = \frac{e^{inx} - e^{-inx}}{2i}, \quad \cos(nx) = \frac{e^{inx} + e^{-inx}}{2}.$$

Let $f \in C^0(\mathbb{R})$ be periodic. The full Fourier series can be written in complex form

$$\sum_{n=-\infty}^{\infty} c_n e^{in\theta} = \mathcal{F}(f)$$

where

$$c_n = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) e^{-inx} dx.$$

To see this we introduce a *Hermitian inner product*.

$$(f, g) = \int_{-\pi}^{\pi} f(x) \overline{g(x)} dx \text{ for } f, g \in C^0(\mathbb{R}, \mathbb{C}) \text{ periodic}$$

where $g(x) = \operatorname{Re}g(x) + i\operatorname{Im}g(x)$ and $\overline{g(x)} = \operatorname{Re}g(x) - i\operatorname{Im}g(x)$ is the complex conjugate of the complex number $g(x)$.

Note that $C^0(\mathbb{R}, \mathbb{C})$ is complex vector space. $X_n = e^{inx}$, $n \in \mathbb{Z}$, are orthogonal w.r.t. (\cdot, \cdot) and

$$\|e^{inx}\|_2^2 = (e^{inx}, e^{inx}) = \int_{-\pi}^{\pi} e^{inx} e^{-inx} dx = 2\pi.$$

In particular, the general Fourier coefficients take the form

$$c_n = \frac{1}{\|X_n\|} (f, X_n) = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) e^{-inx} dx.$$

We can interpret a 2π periodic function f on \mathbb{R} as a function \hat{f} on the 1D circle $\mathbb{S}^1 \subset \mathbb{R}^2 = \mathbb{C}$:

$$f(x) = \hat{f}(e^{ix}).$$

We interpret the heat equation

$$\begin{aligned}u_t &= ku_{x,x} \quad \text{on } \mathbb{R} \\u(x, 0) &= f(x) \quad \text{on } \mathbb{R} \\u(x, t) &= u(x + 2\pi, t) \quad \forall x \in \mathbb{R}\end{aligned}$$

with $f \in C^1(\mathbb{R})$ that is 2π periodic as heat equation on \mathbb{S}^1 .

The solution is given by

$$u(x, t) = \sum_{n=-\infty}^{\infty} c_n e^{-ktn^2} e^{inx} = \sum_{n=-\infty}^{\infty} \int_{-\pi}^{\pi} f(y) \sum_{n=-\infty}^{\infty} e^{-ktn^2} e^{in(x-y)} dy$$

where c_n are the complex Fourier coefficients of f .

We can write this formula as

$$u(x, t) = \int_{-\pi}^{\pi} f(y) K^{\mathbb{S}^1}(x - y, t) dy$$

with

$$K^{\mathbb{S}^1}(\theta, t) = \sum_{n=-\infty}^{\infty} e^{-ktn^2} e^{in\theta}$$

$K^{\mathbb{S}^1}(\theta, t)$ is called the fundamental solution for the heat equation or heat kernel on \mathbb{S}^1 .