

# Laguerre polynomials and Perron-Frobenius operators

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## 1 Laguerre polynomials

### 1.1 Definition and generating functions

Let  $D = \frac{d}{dx}$ . For  $\alpha > -1$  and  $n \geq 0$  let

$$L_n^\alpha(x) = e^x \frac{x^{-\alpha}}{n!} D^n(e^{-x} x^{n+\alpha}),$$

called the **Laguerre polynomials**. Using the Leibniz rule for  $D^n(f \cdot g)$  yields

$$L_n^\alpha(x) = \sum_{k=0}^n \frac{\Gamma(n + \alpha + 1)}{\Gamma(k + \alpha + 1)} \frac{(-x)^k}{k!(n-k)!}.$$

The generating function for the Laguerre polynomials is<sup>1</sup>

$$w(x, z) = (1 - z)^{-\alpha-1} e^{-xz/(1-z)} = \sum_{n=0}^{\infty} L_n^\alpha(x) z^n, \quad |z| < 1.$$

Define

$$W(x, y, z) = (1 - z)^{-1} e^{-(x+y)z/(1-z)} (xyz)^{-\alpha/2} I_\alpha \left( \frac{2\sqrt{xyz}}{1-z} \right), \quad |z| < 1,$$

where

$$I_\alpha(x) = i^{-\alpha} J_\alpha(ix) = \sum_{m=0}^{\infty} \frac{1}{m! \Gamma(m + \alpha + 1)} (x/2)^{2m+\alpha}$$

and

$$J_\alpha(x) = \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m + \alpha + 1)} (x/2)^{2m+\alpha}.$$

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<sup>1</sup>N. N. Lebedev, *Special Functions and Their Applications*, p. 77, §4.17.

$W$  satisfies

$$W(x, y, z) = \sum_{n=0}^{\infty} \frac{n! L_n^\alpha(x) L_n^\alpha(y)}{\Gamma(n + \alpha + 1)} z^n.$$

## 1.2 Differential equations satisfied by Laguerre polynomials

$w$  satisfies the ordinary differential equation

$$(1 - z^2)\partial_z w + (x - (1 - z)(1 + \alpha))w = 0.$$

This yields, for  $n \geq 1$ ,

$$(n + 1)L_{n+1}^\alpha(x) + (x - \alpha - 2n - 1)L_n^\alpha(x) + (n + \alpha)L_{n-1}^\alpha(x) = 0. \quad (1)$$

$w$  also satisfies the ordinary differential equation

$$(1 - t)\partial_x w + tw = 0,$$

which yields, for  $n \geq 1$ ,

$$DL_n^\alpha - DL_{n-1}^\alpha + L_{n-1}^\alpha = 0. \quad (2)$$

Using (1) and (2) gives

$$xDL_n^\alpha = nL_n^\alpha - (n + \alpha)L_{n-1}^\alpha, \quad n \geq 1. \quad (3)$$

Using (3) and (2) we get, for  $n \geq 0$ ,

$$xD^2L_n^\alpha(x) + (\alpha + 1 - x)DL_n^\alpha(x) + nL_n^\alpha(x) = 0. \quad (4)$$

## 1.3 Integral formulas for Laguerre polynomials

For  $\nu > -1$ ,  $a > 0$ ,  $b > 0$ , using the series for  $J_\nu$  one calculates<sup>2</sup>

$$\int_0^\infty e^{-a^2 x^2} J_\nu(bx) x^{\nu+1} dx = \frac{b^\nu}{(2a^2)^{\nu+1}} e^{-\frac{b^2}{4a^2}}. \quad (5)$$

Applying this with  $\nu = n + \alpha$ ,  $a = 1$ ,  $b = 2\sqrt{x}$ ,  $x = \sqrt{t}$  yields

$$\int_0^\infty e^{-t} J_{n+\alpha}(2\sqrt{xt})(\sqrt{t})^{n+\alpha+1} \cdot \frac{1}{2\sqrt{t}} dt = \frac{(2\sqrt{x})^{n+\alpha}}{2^{n+\alpha+1}} e^{-x},$$

i.e.

$$\int_0^\infty e^{-t} J_{n+\alpha}(2\sqrt{xt})(\sqrt{xt})^{n+\alpha} dt = e^{-x} x^{n+\alpha}. \quad (6)$$

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<sup>2</sup>N. N. Lebedev, *Special Functions and Their Applications*, p. 132, §5.15, Example 2.

Now, it is a fact that

$$\frac{d}{du} u^{\nu/2} J_{\nu}(2\sqrt{u}) = u^{(\nu-1)/2} J_{\nu-1}(2\sqrt{u}),$$

and using this and (6), we get that for  $\alpha > 1$  and  $n \geq 0$ ,

$$L_n^{\alpha}(x) = \frac{e^x x^{-\alpha/2}}{n!} \int_0^{\infty} t^{n+\alpha/2} J_{\alpha}(2\sqrt{xt}) e^{-t} dt. \quad (7)$$

We remind ourselves that for  $\alpha > -1$  and  $|z| < 1$ ,

$$(1-z)^{-\alpha-1} e^{-yt/(1-t)} = \sum_{n=0}^{\infty} L_n^{\alpha}(y) z^n.$$

For  $|z| < \frac{1}{3}$ , using this and  $e^{-\frac{yt}{1-t} - \frac{y}{2}} = e^{-\frac{2yt+y-yt}{2(1-t)}} = e^{-\frac{y(1+t)}{2(1-t)}}$  one checks that

$$\begin{aligned} & (1-z)^{-\alpha-1} \int_0^{\infty} e^{-\frac{y(1+t)}{2(1-t)}} y^{\alpha/2} J_{\alpha}(\sqrt{xy}) dy \\ &= \sum_{n=0}^{\infty} z^n \int_0^{\infty} e^{-y/2} y^{\alpha/2} J_{\alpha}(\sqrt{xy}) L_n^{\alpha}(y) dy. \end{aligned}$$

Then one gets, for  $|z| < 1$ ,

$$2e^{-x/2} x^{\alpha/2} \sum_{n=0}^{\infty} (-1)^n L_n^{\alpha}(x) z^n = \sum_{n=0}^{\infty} z^n \int_0^{\infty} e^{-y/2} y^{\alpha/2} J_{\alpha}(\sqrt{xy}) L_n^{\alpha}(y) dy.$$

Therefore for  $\alpha > -1$  and  $n \geq 0$ ,

$$e^{-x/2} x^{\alpha/2} L_n^{\alpha}(x) = \frac{(-1)^n}{2} \int_0^{\infty} J_{\alpha}(\sqrt{xy}) e^{-y/2} y^{\alpha/2} L_n^{\alpha}(y) dy. \quad (8)$$

## 1.4 Orthogonality of Laguerre polynomials

Let

$$\rho_{\alpha}(x) = e^{-x} x^{\alpha}.$$

Let

$$u_n = \rho_{\alpha}^{1/2} L_n^{\alpha}, \quad n \geq 0.$$

$u_n$  satisfies the differential equation

$$(x u_n')' + \left( n + \frac{\alpha+1}{2} - \frac{x}{4} - \frac{\alpha^2}{4x} \right) u_n = 0.$$

Using this we get

$$x(u_n' u_m - u_m' u_n) \Big|_0^{\infty} + (n-m) \int_0^{\infty} u_m u_n dx = 0.$$

Then

$$(n-m) \int_0^\infty u_m u_n dx = 0. \quad (9)$$

Using (1) yields for  $n \geq 2$ ,

$$n(L_n^\alpha)^2 - (n+\alpha)(L_{n-1}^\alpha)^2 - (n+1)L_{n+1}^\alpha L_{n-1}^\alpha + 2L_n^\alpha L_{n-1}^\alpha + (n+\alpha-1)L_n^\alpha L_{n-2}^\alpha = 0.$$

Using this and (9), for  $n \geq 2$ ,

$$n \int_0^\infty e^{-x} x^\alpha L_n^\alpha(x)^2 dx = (n+\alpha) \int_0^\infty e^{-x} x^\alpha L_{n-1}^\alpha(x)^2 dx.$$

Iterating this, for  $n \geq 2$ ,

$$\begin{aligned} \int_0^\infty e^{-x} x^\alpha L_n^\alpha(x)^2 dx &= \frac{(n+\alpha)(n+\alpha-1)\cdots(\alpha+2)}{n(n-2)\cdots 3 \cdot 2} \int_0^\infty e^{-x} x^\alpha L_1^\alpha(x)^2 dx \\ &= \frac{\Gamma(n+\alpha+1)}{n!}. \end{aligned}$$

## 1.5 Asymptotics for Laguerre polynomials

It can be proved that for  $\alpha > -1$ , with  $N = n + \frac{\alpha+1}{2}$ ,<sup>3</sup> for  $x \in \mathbb{R}_{\geq 0}$ ,

$$L_n^\alpha(x) \sim \frac{\Gamma(n+\alpha+1)}{n!} e^{x/2} (Nx)^{-\alpha/2} J_\alpha(2\sqrt{Nx}), \quad n \rightarrow \infty.$$

## 1.6 Laguerre expansions

Suppose that  $f : \mathbb{R}_{>0} \rightarrow \mathbb{R}$  is piecewise smooth in every interval  $[x_1, x_2]$ ,  $0 < x_1 < x_2 < \infty$ , and  $f \in L^2(d\rho_\alpha)$ . Let

$$c_n(f) = \frac{n!}{\Gamma(n+\alpha+1)} \int_0^\infty f(x) L_n^\alpha(x) \rho_\alpha(x) dx,$$

$\rho_\alpha(x) = e^{-x} x^\alpha$ . It can be proved that<sup>4</sup> if  $f$  is continuous at  $x$  then

$$\sum_{n=0}^N c_n(f) L_n^\alpha(x) \rightarrow f(x), \quad N \rightarrow \infty,$$

and if  $f$  is not continuous at  $x$  then

$$\sum_{n=0}^N c_n(f) L_n^\alpha(x) \rightarrow \frac{f(x+0)}{2} + \frac{f(x-0)}{2}, \quad N \rightarrow \infty,$$

which makes sense because  $f$  is a priori piecewise continuous.

<sup>3</sup>N. N. Lebedev, *Special Functions and Their Applications*, p. 87, §4.22.

<sup>4</sup>N. N. Lebedev, *Special Functions and Their Applications*, p. 88, §4.23, Theorem 3.

Let  $\nu > -\frac{1}{2}(\alpha + 1)$  and  $f(x) = x^\nu$ . Integrating by parts,

$$\begin{aligned} c_n(f) &= \frac{n!}{\Gamma(n + \alpha + 1)} \int_0^\infty x^{\nu+\alpha} L_n^\alpha(x) e^{-x} \\ &= \frac{1}{\Gamma(n + \alpha + 1)} \int_0^\infty x^\nu D^n(e^{-x} x^{n+\alpha}) dx \\ &= (-1)^n \frac{\Gamma(\nu + \alpha + 1) \Gamma(\nu + 1)}{\Gamma(n + \alpha + 1) \Gamma(\nu - n + 1)}. \end{aligned}$$

Thus

$$x^\nu = \Gamma(\nu + \alpha + 1) \Gamma(\nu + 1) \sum_{n=0}^{\infty} \frac{(-1)^n L_n^\alpha(x)}{\Gamma(n + \alpha + 1) \Gamma(\nu - n + 1)}.$$

For  $p$  a positive integer,

$$x^p = \Gamma(p + \alpha + 1) \cdot p! \sum_{n=0}^p \frac{(-1)^n L_n^\alpha(x)}{\Gamma(n + \alpha + 1) \cdot (p - n)!}.$$

Define

$$f(x) = (ax)^{-\alpha/2} J_\alpha(2\sqrt{ax}), \quad \alpha > -1, \quad a > 0, \quad x > 0.$$

Using

$$(1 - z)^{-\alpha-1} e^{-xz/(1-z)} = \sum_{n=0}^{\infty} L_n^\alpha(x) z^n, \quad |z| < 1,$$

we obtain, as  $e^{-\frac{xz}{1-z}} = e^{-x/(1-z)}$ ,

$$\begin{aligned} &(1 - z)^{-\alpha-1} \int_0^\infty e^{-x/(1-z)} (x/a)^{\alpha/2} J_\alpha(2\sqrt{ax}) dx \\ &= \int_0^\infty e^{-x} (x/a)^{\alpha/2} J_\alpha(2\sqrt{ax}) \sum_{n=0}^{\infty} L_n^\alpha(x) z^n dx \\ &= \sum_{n=0}^{\infty} \left( \int_0^\infty f(x) L_n^\alpha(x) \rho_\alpha(x) dx \right) z^n. \end{aligned}$$

Doing the change of variable  $2\sqrt{ax} = by$  with  $b > 0$  and then applying (6) with

$$A^2 = \frac{b^2}{4a(1-z)} \text{ and } \nu = \alpha,$$

$$\begin{aligned} & (1-z)^{-\alpha-1} \int_0^\infty e^{-x/(1-z)} (x/a)^{\alpha/2} J_\alpha(2\sqrt{ax}) dx \\ &= (1-z)^{-\alpha-1} (2a)^{-\alpha-1} b^{\alpha+2} \int_0^\infty e^{-\frac{b^2 y^2}{4a(1-z)}} J_\alpha(by) y^{\alpha+1} dy \\ &= (1-z)^{-\alpha-1} (2a)^{-\alpha-1} b^{\alpha+2} \cdot \frac{b^\alpha}{(2A^2)^{\alpha+1}} e^{-\frac{b^2}{4A^2}} \\ &= (1-z)^{-\alpha-1} (2a)^{-\alpha-1} b^{\alpha+2} \cdot b^{-\alpha-2} (2a(1-z))^{\alpha+1} e^{-a(1-z)} \\ &= e^{-a(1-z)} \\ &= e^{-a} \sum_{n=0}^\infty \frac{(az)^n}{n!}. \end{aligned}$$

Therefore

$$e^{-a} \sum_{n=0}^\infty \frac{a^n}{n!} z^n = \sum_{n=0}^\infty \left( \int_0^\infty f(x) L_n^\alpha(x) \rho_\alpha(x) dx \right) z^n,$$

whence, for  $n \geq 0$ ,

$$c_n(f) = \frac{n!}{\Gamma(n+\alpha+1)} \int_0^\infty f(x) L_n^\alpha(x) \rho_\alpha(x) dx = \frac{n!}{\Gamma(n+\alpha+1)} e^{-a} \frac{a^n}{n!}.$$

Therefore, for  $\alpha > -1$ ,  $a > 0$ ,  $x > 0$ ,

$$(ax)^{-\alpha/2} J_\alpha(2\sqrt{ax}) = \sum_{n=0}^\infty c_n(f) L_n^\alpha(x) = e^{-a} \sum_{n=0}^\infty \frac{a^n}{\Gamma(n+\alpha+1)} L_n^\alpha(x).$$

## 2 Integral operators

We remind ourselves that, for  $\alpha = 1$ ,

$$u_n(x) = \rho_1(x)^{1/2} L_n^1(x) = e^{-x/2} x^{1/2} L_n^1(x).$$

$\{u_n : n \geq 0\}$  is an orthonormal basis for  $L^2(\mathbb{R}_{\geq 0})$ .

For  $x, y \in \mathbb{R}_{> 0}$  define

$$k(x, y) = k_x(y) = k^x(y) = \frac{J_1(2\sqrt{xy})}{((e^x - 1)(e^y - 1))^{1/2}}.$$

For  $\phi \in L^2(\mathbb{R}_{\geq 0})$  and  $y \in \mathbb{R}_{> 0}$ , define

$$K\phi(y) = \int_{\mathbb{R}_{\geq 0}} k_y(x) \phi(x) dx.$$

We have established, with  $\alpha = 1$ ,

$$J_1(2\sqrt{xy}) = (xy)^{1/2}e^{-x} \sum_{n=0}^{\infty} \frac{x^n}{(n+1)!} L_n^1(y).$$

Hence

$$\begin{aligned} \int_0^{\infty} k_y(x)\phi(x)dx &= \int_0^{\infty} \phi(x)(e^x - 1)^{-1/2}(e^y - 1)^{-1/2}(xy)^{1/2}e^{-x} \\ &\quad \cdot \sum_{n=0}^{\infty} \frac{x^n}{(n+1)!} L_n^1(y) dx \\ &= \sum_{n=0}^{\infty} \frac{(e^y - 1)^{-1/2}y^{1/2}L_n^1(y)}{(n+1)!} \int_0^{\infty} \phi(x)(e^x - 1)^{-1/2}x^{1/2}e^{-x}x^n dx \\ &= \sum_{n=0}^{\infty} q_n(y) \langle \phi, p_n \rangle, \end{aligned}$$

for

$$p_n(x) = \frac{1}{(n+1)!} (e^x - 1)^{-1/2} e^{-x} x^{n+\frac{1}{2}} = \frac{1}{(n+1)!} e^{-x/2} (e^x - 1)^{-1/2} x^n u_n(x)$$

and

$$q_n(y) = (e^y - 1)^{-1/2} y^{1/2} L_n^1(y) = (1 - e^{-y})^{-1/2} u_n(y).$$

Then

$$K\phi = \sum_{n=0}^{\infty} q_n \langle \phi, p_n \rangle.$$

The following states the trace of the operator  $K : L^2(\mathbb{R}_{\geq 0}) \rightarrow L^2(\mathbb{R}_{\geq 0})$ .<sup>5</sup>

**Theorem 1.**  $\text{tr } K = \int_0^{\infty} k(x, x) dx = \int_0^{\infty} \frac{J_1(2x)}{(e^x - 1)} dx = 0.7711 \dots$

### 3 Hardy spaces

For  $x \in \mathbb{R}$  let  $P_x = \{z \in \mathbb{C} : \text{Re } z > x\}$ . Let  $H$  be the collection of holomorphic functions  $f : P_{-1/2} \rightarrow \mathbb{C}$  such that for any  $x > -\frac{1}{2}$ ,  $f|_{P_x}$  is bounded and such that

$$\int_{\mathbb{R}} \left| f\left(-\frac{1}{2} + iy\right) \right|^2 dy < \infty.$$

Define  $M : L^2(\mathbb{R}_{\geq 0}) \rightarrow H$ , for  $\phi \in L^2(\mathbb{R}_{\geq 0})$ , by

$$M\phi(z) = \int_{\mathbb{R}_{\geq 0}} e^{-zs-s/2} \phi(s) ds.$$

<sup>5</sup>cf. A. A. Kirillov, *Elements of the Theory of Representations*, p. 211, §13, Theorem 2.

For  $f \in H$  define

$$P_\lambda f(z) = \sum_{k \geq 1} \frac{1}{(z+k)^2} f\left(\frac{1}{z+k}\right) \quad \operatorname{Re} z > -\frac{1}{2},$$

called a **Perron-Frobenius operator**.  $\lambda$  denotes Lebesgue measure.

Let

$$h(s) = \left(\frac{1-e^{-s}}{s}\right)^{1/2}$$

for  $s \in \mathbb{R}_{>0}$ , with  $h(0) = 1$ . Because  $h \in L^\infty(\mu)$ , it makes sense to define  $S : L^2(\mathbb{R}_{\geq 0}) \rightarrow L^2(\mathbb{R}_{\geq 0})$  by

$$S\phi(s) = h\phi, \quad \phi \in L^2(\mathbb{R}_{\geq 0}).$$

Define  $A : H \rightarrow L^2(\mathbb{R}_{\geq 0})$  by

$$A = S \circ M^{-1}.$$

We prove that  $P_\lambda$  and  $K$  are conjugate.<sup>6</sup>

**Theorem 2.**  $P_\lambda = A^{-1}KA$ .

*Proof.* Let  $\phi \in L^2(\mathbb{R}_{\geq 0})$  and set  $f = M\phi$ . Then

$$A^{-1}KAf = A^{-1}KS\phi.$$

We calculate

$$\begin{aligned} (S^{-1}KS\phi)(x) &= h(x)^{-1} \int_{\mathbb{R}_{\geq 0}} k_x(y) \cdot h(y) \cdot \phi(y) dy \\ &= \left(\frac{x}{1-e^{-x}}\right)^{1/2} \int_0^\infty \frac{J_1(2\sqrt{xy})}{((e^x-1)(e^y-1))^{1/2}} \cdot \left(\frac{1-e^{-y}}{y}\right)^{1/2} \cdot \phi(y) dy \\ &= \int_0^\infty \left(\frac{x}{y}\right)^{1/2} \frac{e^{x/2}}{(e^x-1)^{1/2}} \frac{(e^y-1)^{1/2}}{e^{y/2}} \frac{J_1(2\sqrt{xy})}{((e^x-1)(e^y-1))^{1/2}} \phi(y) dy \\ &= \int_0^\infty \left(\frac{x}{y}\right)^{1/2} \frac{e^{(x-y)/2}}{e^x-1} J_1(2\sqrt{xy}) \phi(y) dy. \end{aligned}$$

Then

$$\begin{aligned} &(MS^{-1}KS\phi)(z) \\ &= \int_{\mathbb{R}_{\geq 0}} e^{-zx-x/2} (S^{-1}KS\phi)(x) dx \\ &= \int_0^\infty e^{-zx-x/2} \left( \left(\frac{x}{y}\right)^{1/2} \frac{e^{(x-y)/2}}{e^x-1} J_1(2\sqrt{xy}) \phi(y) dy \right) dx \\ &= \end{aligned}$$

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<sup>6</sup>Marius Iosifescu and Cor Kraaikamp, *Metrical Theory of Continued Fractions*, p. 9, Proposition 1.1.1.



It is a fact that for  $\operatorname{Re} z > -1$  and for  $t \geq 0$ ,

$$\sum_{k \geq 0} (z+k)^{-2} \exp\left(-\frac{t}{z+k}\right) = \int_0^\infty (st^{-1})^{1/2} e^{-zs} \frac{J_1(2\sqrt{st})}{e^s - 1} ds.$$

Using this,

$$\begin{aligned} (MS^{-1}KS\phi)(z) &= \int_0^\infty e^{-y/2} \left( \int_0^\infty (xy^{-1})^{1/2} e^{-zx} \frac{J_1(2\sqrt{xy})}{e^x - 1} dx \right) \phi(y) dy \\ &= \int_0^\infty e^{-y/2} \sum_{k \geq 1} (z+k)^{-2} \exp\left(-\frac{y}{z+k}\right) \cdot \phi(y) dy \\ &= \sum_{k \geq 1} (z+k)^{-2} \left( \int_0^\infty \exp\left(-\frac{y}{z+k} - \frac{y}{2}\right) \phi(y) dy \right) \\ &= \sum_{k \geq 1} (z+k)^{-2} \cdot M\phi\left(\frac{1}{z+k}\right). \end{aligned}$$

Thus, as  $f = M\phi$ ,

$$(MS^{-1}KSM^{-1}f)(z) = \sum_{k \geq 1} (z+k)^{-2} f\left(\frac{1}{z+k}\right) = P_\lambda f(z),$$

that is,

$$A^{-1}KAf(z) = P_\lambda f(z).$$

□